

Structural Engineering Three-Dimensional Nano-Heterojunction Networks for High-Performance Photochemical Sensing

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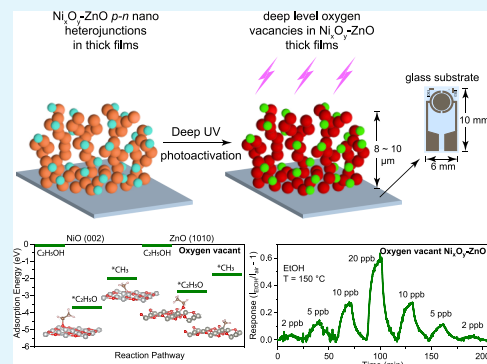
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ABSTRACT: Nanoscale heterojunction networks are increasingly regarded as promising functional materials for a variety of optoelectronic and photocatalytic devices. Despite their superior charge-carrier separation efficiency, a major challenge remains in the optimization of their surface properties, with surface defects playing a major role in charge trapping and recombination. Here, we report the effective engineering of the photocatalytic properties of nanoscale heterojunction networks *via* deep ultraviolet photoactivation throughout their cross-section. For the first time, in-depth XPS analysis of very thick ($\sim 10 \mu\text{m}$) $\text{Ni}_x\text{O}_y\text{-ZnO}$ films reveals localized *p-n* nanoheterojunctions with tunable oxygen vacancies (V_o) originating from both Ni_xO_y and ZnO nanocrystals. Optimizing the amount of oxygen vacancies leads to a 30-fold increase in the photochemoresistive response of these networks, enabling the detection of representative analyte concentrations down to 2 and 20 ppb at an optimal temperature of 150°C and room temperature, respectively. Density functional theory calculations reveal that this performance enhancement is presumably due to an 80% increase in the analyte adsorption energy. This flexible nanofabrication approach in conjunction with straightforward vacancy control *via* photoactivation provides an effective strategy for engineering the photocatalytic activity of porous metal oxide semiconductor networks with applications in chemical sensors, photodetectors, and photoelectrochemical cells.

KEYWORDS: nanoscale heterojunctions, oxygen vacancy, room temperature sensing, deep ultraviolet photoactivation, ZnO-NiO, XPS, DFT



INTRODUCTION

Nanoscale heterojunctions have recently emerged as a promising multifunctional material platform. Among other benefits, the engineering of strong electrical field gradients in these nanoscale structures allows to attain nonconventional optical, electrical, and (electro)chemical properties that are not achievable with single semiconductor materials and their bulk heterojunctions.¹ Advances in the physics and fabrication of nanoscale heterojunctions are leading a renaissance in the architecture and composition of optoelectronic devices, demonstrating high-efficiency solar cells,² light-emitting diodes,³ transistors,⁴ biomolecular,⁵ and chemical sensors⁶ as well as potential for quantum computing and spintronics.⁷ While significant progress has been made in the understanding of their electronic structure and energy level alignment,⁸ the large specific surface of nanoscale heterojunctions is prone to induce trapping states and high recombination rates that may negatively impact device performance. Therefore, the judicious engineering of the surface composition of nanoscale heterojunctions is highly desired to realize their full potential in photochemical and optoelectronic applications.

The use of nanoscale heterojunctions for sensing has been pioneered by the engineering of one-⁹ and zero-dimensional¹⁰

ZnO-NiO nanocomposite structures for visible-blind UV photodetectors, demonstrating vastly superior device performance such as enhanced detectivity, rise and recovery times. These findings have been translated into a variety of photocatalytic and optoelectronic applications, including the development of high-surface-area ZnO-NiO particle networks for low-temperature photochemical sensing.¹¹ The formation of $\text{NiO-Fe}_2\text{O}_3$ *p-n* nanoscale heterojunctions has also been reported to increase the chemoresistive response of $\alpha\text{-Fe}_2\text{O}_3$ to toluene gas molecules by 7.8 times.¹² This was attributed to more efficient charge separation¹⁰ and enhanced surface reaction kinetics. Similarly, electron–hole separation was reported at *p-CuO/n-TiO*₂ heterointerfaces, resulting in larger depletion zones and more dissociation of analytes.¹³ Recently, Li *et al.*¹⁴ reported an enhanced photocatalytic activity of ZnO-NiO heterojunctions for dye degradation, further indicating a

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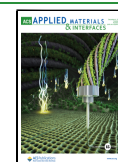


Table 1. Comparison of Recently Developed Chemical Sensors Based on Heterostructures with Oxygen Defects

Material	Sensing temperature (°C)	Oxygen vacancy introduction method	EtOH concentration (ppm)	Responsivity ($I_{\text{EtOH}}/I_{\text{air}} - 1$)	Detection limit (ppm)	ref
Porous Ni _x O _y -ZnO nanoscale heterojunction networks	RT ^a	Photoactivation at 200 °C	0.05	1.91	0.01	This work
	150	Photoactivation at 200 °C	1	8.38		
ZnO-NiO nano heterojunction	RT ^a	Without oxygen vacancies	1	2.6	0.01	11
Co-doped ZnO microspheres	220	Co-doping at 400 °C	5	3.3 ^b	N/A	25
CeO ₂ -doped ZnO nanostructures	350	UV irradiation method, calcination at 600 °C	500	150 ^b	N/A	26
Ce-doped ZnO nanostructures	300	Ce-doping at 450 °C	100	72.6 ^b	N/A	27
ZnO/SnO ₂ composite hollow spheres	225	Hydrothermal process, calcination at 400 °C	30	34.8 ^b	0.5	28
ZnO-In ₂ O ₃ nanostructures	250	Solution synthesis, calcination at 400 °C	300	900 ^b	N/A	29
ZnO added MoO ₃ nanostructures	300	Sputtering, annealing at 500 °C	300	38 ^b	N/A	30
Tb-SnO ₂ nanotubes	200	Electrospinning process, calcination at 500 °C	100	53.6 ^b	1	31
CdO-loaded ZnO nanostructures	250	Hydrothermal process, annealing at 500 °C	100	65.5 ^b	0.5	32
CeO ₂ decorated SnO ₂ hollow spheres	225	Hydrothermal process, calcination at 400 °C	100	37 ^b	10	33

^aRT = room temperature. ^b $I_{\text{EtOH}}/I_{\text{air}}$.

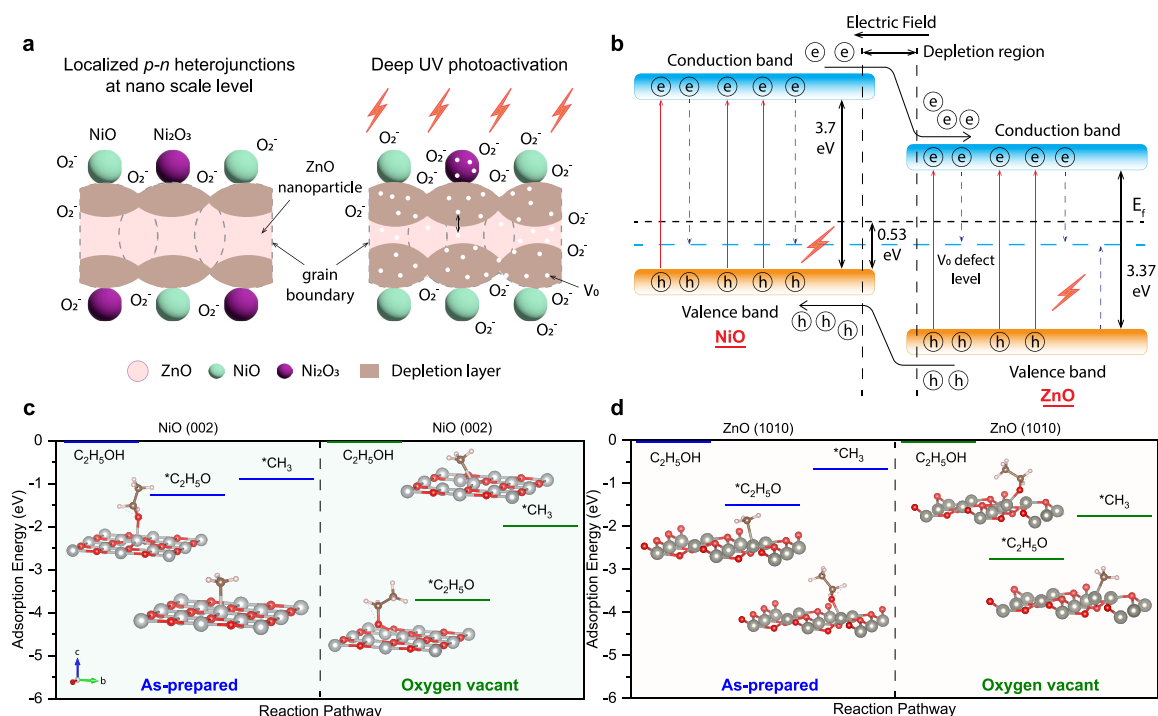


Figure 1. Schematic illustration of (a) *p-n* nanoscale heterojunctions formed by the Ni_xO_y-ZnO nanoparticles and (b) their energy band diagram upon induction of the oxygen vacancies, showing the formation of an internal electric field at the nanoscale heterointerface with the vector pointing from positively charged *n*-ZnO to negatively charged *p*-Ni_xO_y. DFT calculations with corresponding geometric configurations of ethanol dissociation and adsorption on (c) as-prepared (without DUV photoactivation) and oxygen vacant (with DUV photoactivation) *p*-NiO (green shaded) and (d) *n*-ZnO (yellow shaded) surfaces with and without DUV photoactivation. The reaction pathway shows the physisorption of the ethanol molecule followed by its dissociation (chemisorption) into C₂H₅O and CH₃ by redox reactions on both surfaces.

potential role of the nanocomposite surface structure in enhancing surface interactions.

In parallel, semiconductor defect engineering, and in particular, induction of oxygen vacancies, has been explored to enhance the efficiency of various photocatalytic and photoelectrochemical processes,^{15–17} including photoelectrochemical hydrogen production,¹⁸ and organic photodecomposition.¹⁹ Introducing oxygen vacancies in a semiconductor can

enhance the surface density of redox sites, increasing catalytic activity.^{15,19,20} Zhang *et al.* optimized the chemical and electronic properties of NiO nanorods for electrocatalysis by oxygen vacancy engineering,²¹ significantly enhancing the electrical conductivity and hydrogen evolution reaction kinetics. Li *et al.*²² found that surface oxygen vacancies in ZnO nanofilms enhanced their photochemical sensing response to NO₂ gas molecules, with neutral oxygen and

doubly ionized oxygen vacancies favoring the adsorption of O₂ and NO₂ adsorption, respectively. Recently, defect engineering has been applied to enhance the electronic and photocatalytic properties of nanoscale heterojunctions, improving light harvesting,¹⁵ charge separation and transfer,²³ and surface reaction kinetics.^{15,24}

Despite this potential, current approaches for the induction of oxygen vacancies in nanoscale heterojunctions have significant limitations. Oxygen vacancies are often generated in heterostructures by doping, which requires relatively high temperatures and carefully controlled reaction environments (Table 1).²¹ Most approaches also require postfabrication annealing. This makes it challenging to control the concentration of oxygen vacancies and provides limitations for the use of low-melting-temperature substrates. For instance, Li *et al.* introduced surface oxygen vacancies in ZnO nanofilms by annealing between 200 and 450 °C under UV light.²² Xu *et al.* generated surface oxygen vacancies by Co doping of ZnO microspheres at 400 °C.²⁵ Hamedani *et al.* induced surface oxygen vacancies by doping CeO₂ in ZnO nanostructures under ultraviolet irradiation at 600 °C.²⁶ Furthermore, there is a lack of studies and approaches for the engineering of oxygen vacancies in thick porous films required to provide suitable surface area for many (photo)-electrocatalytic applications. Recently, Zhang *et al.* reported the ammonia sensing performance of dense films of oxygen-deficient ZnO-NiO nanoscale heterojunctions.²⁴ However, the low film porosity resulted in a relatively poor chemical sensing performance, with an ammonia limit of detection of 3000 ppm.

Here, we introduce a flexible approach for the engineering and fabrication of thick nanoscale heterojunction films with tailored surface and structural properties and excellent photochemical activity. Ultraporos nanoscale heterojunction networks of electron-depleted Ni_xO_y-ZnO nanoparticles were self-assembled by a scalable gas-phase approach. Oxygen defects were engineered in these Ni_xO_y-ZnO networks with a thickness of up to 10 μm *via* mild-temperature deep ultraviolet (DUV) photoactivation, as confirmed by X-ray photoelectron spectroscopy depth profiling. The induction of oxygen vacancies on the surface and throughout the nanoscale heterojunctions drastically increased their photocatalytic performance, resulting in a 30-fold enhancement of their photochemical sensing response to a representative gas analyte and a record-low limit of detection of 20 ppb at room temperature. The latter enhancement was explained by density functional theory (DFT) calculations, revealing an 80% increase in the analyte gas molecule adsorption energy upon oxygen vacancy induction. These insights provide an effective method to optimize the photocatalytic and optoelectronic properties of high-surface-area nanoscale heterojunction films with applications in a variety of emerging optoelectronic and photocatalytic devices.

RESULTS AND DISCUSSION

Figure 1a shows a schematic description of our approach for engineering the surface properties of thick nanoscale heterojunction films *via* DUV photoactivation under inert ambient conditions. Studies of surface oxygen vacancies in nanoscale heterojunction networks suggest that the formation of *p-n* nanoscale heterojunctions can increase surface band bending resulting in a larger surface depletion region,³⁴ which is beneficial to enhance the electrical transduction of both photo-³⁴ and chemical³⁵ interaction events. Furthermore, the

formation of localized midlevel electronic states, due to the presence of oxygen vacancies, can further increase the photoexcited charge separation efficiency, adding to the electrical field-based separation mechanism of nanoscale heterojunctions (Figure 1b).

We performed DFT modeling of the interaction of a representative analyte, ethanol, with the surface of pristine and oxygen vacant Ni_xO_y and ZnO surfaces (see Experimental Section). Figure 1c, d show that ethanol adsorbs on the pristine (as-prepared) and oxygen vacant Ni_xO_y and ZnO surfaces, respectively. The initial adsorption of ethanol on the surface is physisorption where ethanol molecules make contact with the metal oxide surfaces. Ethanol molecules adsorb on the ZnO surface and dissociate into intermediate components. There are various routes for ethanol disassociation on metal oxide surfaces and the associated charge carrier migration is quite intricate.³⁶ However, the proposed low-temperature physisorption and chemisorption processes are presented in Figure S1b.

After the ethanol adsorption, ethanol and the metal oxides participate in a redox reaction (chemisorption) and the ethanol molecule reduces to C₂H₅O (Figure S1b, Step 2). The C₂H₅O radicals have three isomers, ethoxy (CH₃CH₂O), 1-hydroxyethyl (CH₃CHOH), and 2-hydroxyethyl (CH₂CH₂OH) radicals.³⁷ These radicals play a pivotal role as reaction intermediates in the decomposition and combustion of ethanol. C₂H₅O can further decompose by 3 different reaction pathways; carbon-carbon bond fission, carbon-hydrogen bond fission, and isomerization followed by decomposition.^{37,38} It has been reported in separate studies that the carbon-carbon bond fission (direct decomposition; Figure S1b, reaction path 1) is an order of magnitude faster than that of carbon-hydrogen bond fission and has the lowest reaction barrier as compared to reaction paths 2 and 3.^{37,38} Our DFT calculations show higher adsorption energies of C₂H₅O on as-prepared Ni_xO_y (1.27 eV) and ZnO (1.51 eV), while for CH₃, it is 0.89 and 0.65 eV on the Ni_xO_y and ZnO surfaces, respectively, indicating the dominant decomposition pathway *via* reaction path 1. The induction of oxygen vacancies leads to a significant increase in the adsorption energy of C₂H₅O to 3.7 and 2.78 eV on the Ni_xO_y and ZnO surfaces, respectively. Similarly, the CH₃ adsorption energy increases to 2.0 and 1.78 eV on the oxygen vacant Ni_xO_y and ZnO surfaces, respectively (Figure 1d). This suggests that the induction of oxygen vacancies in the Ni_xO_y-ZnO nanoscale heterojunctions can increase the adsorption of ethanol gas molecules and facilitate their further reaction with the surface.

The higher adsorption energy on the Ni_xO_y is in line with its excellent surface activity,³⁹ making it one of the preferred *p*-type semiconductors for nanoheterojunction-based UV and photochemical sensing applications.^{10,11} Recent models suggest that polycrystalline *p*-NiO has a strong ethanol sensing response at room temperature, due to enhanced electron transfer.⁴⁰ The high sensitivity and fast response to ethanol gas molecules of ZnO (10 $\bar{1}$ 0) nanorods were attributed to the higher adsorption energies of ethanol, NH₃, H₂, CO, and O₂, computed by DFT calculations.⁴¹ The latter was revealed to arise from an adsorption-induced reconstruction of the (10 $\bar{1}$ 0) ZnO surface, which caused changes in the electronic structure. Similarly, Xu *et al.*⁴² utilized DFT to explore the interactions of NO with the (110) surface of SnO₂, showing that the adsorption energy of NO on an oxygen-deficient surface is greater than on a stoichiometric one. It was also found that an

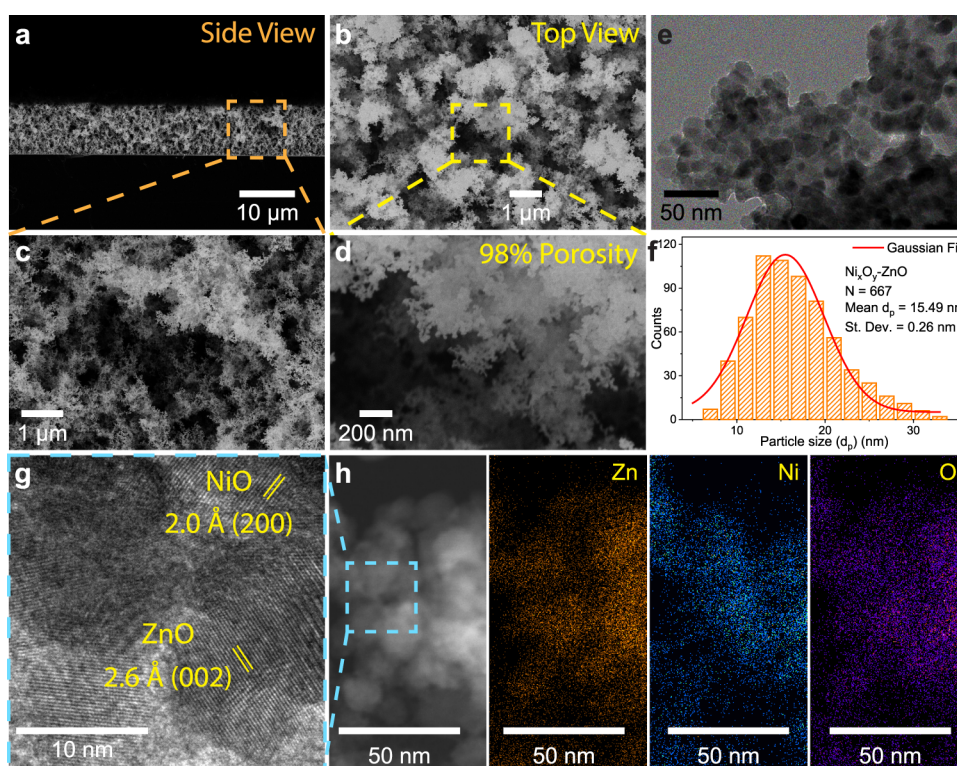


Figure 2. Low magnification scanning electron microscopic (SEM) images of the ultraporous $\text{Ni}_x\text{O}_y\text{-ZnO}$ *p-n* nanoscale heterojunction networks: (a) cross-sectional and (b) top view. Corresponding high-resolution SEM images of (c) cross-sectional and (d) top view revealing high porosity in the nanoscale *p-n* heterojunction fractals. (e) Low-resolution transmission electron microscopic (TEM) image of $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks and (f) $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoparticle size analysis by TEM images showing an average size of 15.5 nm. (g) High-resolution TEM (HRTEM) of $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks showing the *d*-spacing of ZnO (2.6 Å), corresponding to (002) and NiO (2.0 Å), corresponding to (200) plane. (h) Backscattered TEM and the corresponding electron energy-dispersive X-ray spectroscopy (EDS) elemental mapping of $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks.

oxygen-deficient surface with physisorbed oxygen exhibits a further increase in the interaction with NO. The importance of the surface's stoichiometry, including the presence of oxygen vacancies, was further corroborated by other studies.⁴³ Here, we proposed that the oxygen vacancies induced by DUV irradiation improve the charge separation efficiency, increase the depletion region, and enhance charge transfer rates between the analyte molecule and the nanoscale heterojunction networks, thereby promoting the analyte adsorption and dissociation on the metal oxide surfaces and at the nanoscale heterojunction networks with significantly high adsorption energies.

In addition to the surface properties, the photochemiresistive performance of nanoscale heterojunctions depends on the porosity and thickness of the layer. Here, 98% porous networks of ZnO nanoparticles with a thickness of 10 μm were self-assembled by flame spray pyrolysis (FSP) (Figure S1a) and decorated by Ni sputtering and postsintering (see Experimental Section), as previously reported.¹¹ A low magnification cross-sectional scanning electron microscopic (SEM) image of the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks reveals their uniform thickness (8–10 μm) across the substrate, with no cracks or visible defects (Figure 2a). Aerosol deposition allows precise control over the network thickness, porosity, and average nanoparticle size, resulting in highly pure ZnO films.^{10,44} Here, thickness (8–10 μm) and porosity (~98%) were obtained by controlling the flame exposure time and the height above the burner, which were optimized at 100 s and 20 cm, respectively.

For the first time, oxygen vacancies were induced in these nanoscale heterojunction networks by DUV photoactivation at 200 °C using an excimer lamp (30 mW cm^{-2}) in a nitrogen atmosphere. The short-wavelength and high-energy DUV irradiation (172 nm and 7.2 eV) were effective for deep penetration of DUV light into the ultraporous network morphology at a low temperature. This enabled the creation of oxygen vacancies not only on the surface but also several micrometers below the top surface while preserving the integrity of nanoscale heterojunction networks. Figure 2b presents a low magnification top-view SEM image of the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks. The corresponding high magnification SEM images (Figure 2c,d) reveal the presence of 1–4 μm macro- and 500 nm mesopores, in line with metal oxide nanoparticle networks previously synthesized via the flame-aerosol route.^{44,45}

SEM images at various magnifications show a self-similar fractal structure (Figure S2) with a porosity of ca. 98% (Experimental Section). These $\text{Ni}_x\text{O}_y\text{-ZnO}$ fractals were analyzed for their fractal dimensions and lacunarity. The fractal dimension, D_f is a dimensionless number that represents the complexity and self-similarity of a fractal system, usually ranging from 1.75 to 2.5 for aerosol-deposited materials. It can be measured using scattering techniques such as light, X-ray, and neutron scattering, or image analysis.⁴⁵ Figure 2c was used to calculate the D_f and lacunarity of $\text{Ni}_x\text{O}_y\text{-ZnO}$ fractals using image analysis and the fractal image generated by pixel distribution is presented in Figure S2c (see Experimental Section). Here, a fractal dimension of $1.81 \pm$

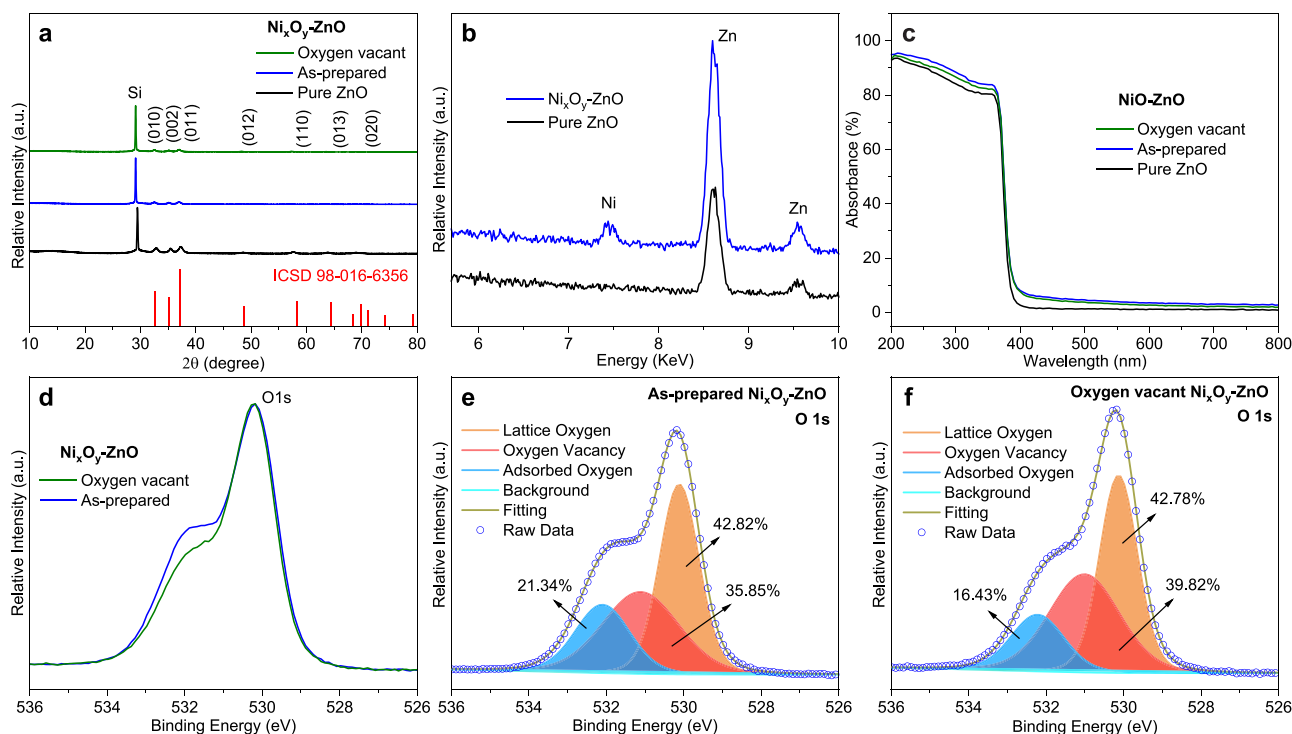


Figure 3. X-ray diffraction (XRD) patterns of (a) Ni_xO_y-ZnO nanoscale heterojunction networks with pure ZnO and the reference patterns. (b) Energy dispersive spectra (EDS) of Ni_xO_y-ZnO nanoscale heterojunction networks. (c) Optical absorbance. (d) High-resolution O 1s X-ray photoelectron spectroscopy (XPS) spectra of Ni_xO_y-ZnO nanoscale heterojunction networks. Deconvoluted O 1s spectra of (e) as-prepared and (f) oxygen vacant Ni_xO_y-ZnO nanoscale heterojunction networks.

0.04 was computed for these Ni_xO_y-ZnO nanoscale heterojunction networks using the image processing and box-counting algorithms (see [Experimental Section](#)), as visualized in the slope of $\ln N$ versus $\ln \epsilon$ plot in [Figure S3a](#). The lacunarity, Λ , quantifies the inhomogeneity of fractal structures and was calculated from the variance, σ , and the mean, μ , of the foreground pixel distribution ([Figure S2c](#)). The plot of the lacunarity versus box length in [Figure S3b](#) shows the maximum between 0 and 0.4 μm (0–400 nm), indicating a higher degree of inhomogeneity and uniform disorder on this length scale. These values are in line with previous reports on flame-aerosol-made films and a layer porosity of 98%.^{45,46}

The compositional and morphological configuration of the localized Ni_xO_y-ZnO nanoscale heterojunction networks were studied using transmission electron microscopy (TEM). The nanoparticles were found to be quasi-spherical ([Figure 2e](#)) with a mean particle diameter, d_p , of 15.5 ± 0.26 nm, estimated by measuring 667 nanoparticles ([Figure 2f](#)) with an estimated specific surface area of ZnO base sensing layer to be $690 \text{ m}^2/\text{g}$ (Text S1). High-resolution TEM (HRTEM) revealed a d -spacing of 2.6 Å corresponding to the (002) plane of ZnO and 2.0 Å corresponding to the (200) plane of NiO ([Figure 2g](#)).²⁴ Backscattered TEM and electron energy-dispersive X-ray spectroscopy (EDS) elemental mapping analyses were also conducted to confirm the presence and morphology of Ni_xO_y on the ZnO nanoparticle networks ([Figure 2h](#)). [Figure 2g–h](#) reveal that the Ni_xO_y nanoparticles were in close contact with the ZnO nanoparticles forming a network of nanoscale heterojunctions. The analysis in [Figure 2h](#) also suggests a uniform distribution of Ni_xO_y in the ZnO nanoparticle network.

X-ray diffraction (XRD) analysis shows a highly crystalline nature and purity of the nanoscale heterojunction networks.

The XRD patterns of pure ZnO, as-prepared Ni_xO_y-ZnO, and oxygen vacant Ni_xO_y-ZnO are shown in [Figure 3a](#). All of the peaks match the reference pattern of zincite (ICSD 98–016–6356) with a hexagonal crystal system and the space group of $P63mc$. No Ni peaks were found in the XRD patterns of both the as-prepared and the oxygen vacant Ni_xO_y-ZnO networks. This was attributed to the low content of Ni as compared to the ZnO nanoparticles. The average crystal sizes for the as-prepared and oxygen vacant Ni_xO_y-ZnO nanoscale heterojunction networks were 11.7 ± 2.1 nm and 11.8 ± 2.0 nm.⁴⁷ These results align with the average d_p of 15.5 nm calculated by TEM analysis, indicating the presence of mostly single-crystal nanoparticles. These results also show that the DUV photoactivation did not affect the grain size of the metal oxide nanoscale heterojunction networks, making them suitable for catalysis and detection applications.

To confirm the composition of the sputtered Ni_xO_y, pure Ni_xO_y films were made with the same conditions and analyzed using Grazing Incidence XRD (GI-XRD) with an incident angle of 5 degrees and step size of 0.052 degrees ([Figure S4a](#)). The peaks matched the reference pattern of Bunsenite (ICSD 98–002–4014) with a cubic crystal system and the $Fm\bar{3}m$ space group. Subsequently, the Ni_xO_y-ZnO nanoscale heterojunction networks were also analyzed using GI-XRD at an incident angle of 3° with a step size of 0.11°. However, as in the case of the standard XRD analysis, the key peak of NiO at $\sim 43^\circ$ was not detected ([Figure S4b](#)). This is due to the finer distribution of the Ni_xO_y nanoparticles in the highly porous ZnO nanoparticle networks, which lead to very small Ni_xO_y crystal sizes.

To further ascertain the presence of Ni_xO_y, EDS analysis was performed for the pure ZnO and Ni_xO_y-ZnO networks. The EDS spectra of Ni_xO_y-ZnO show peaks for both Ni (7.4 keV)

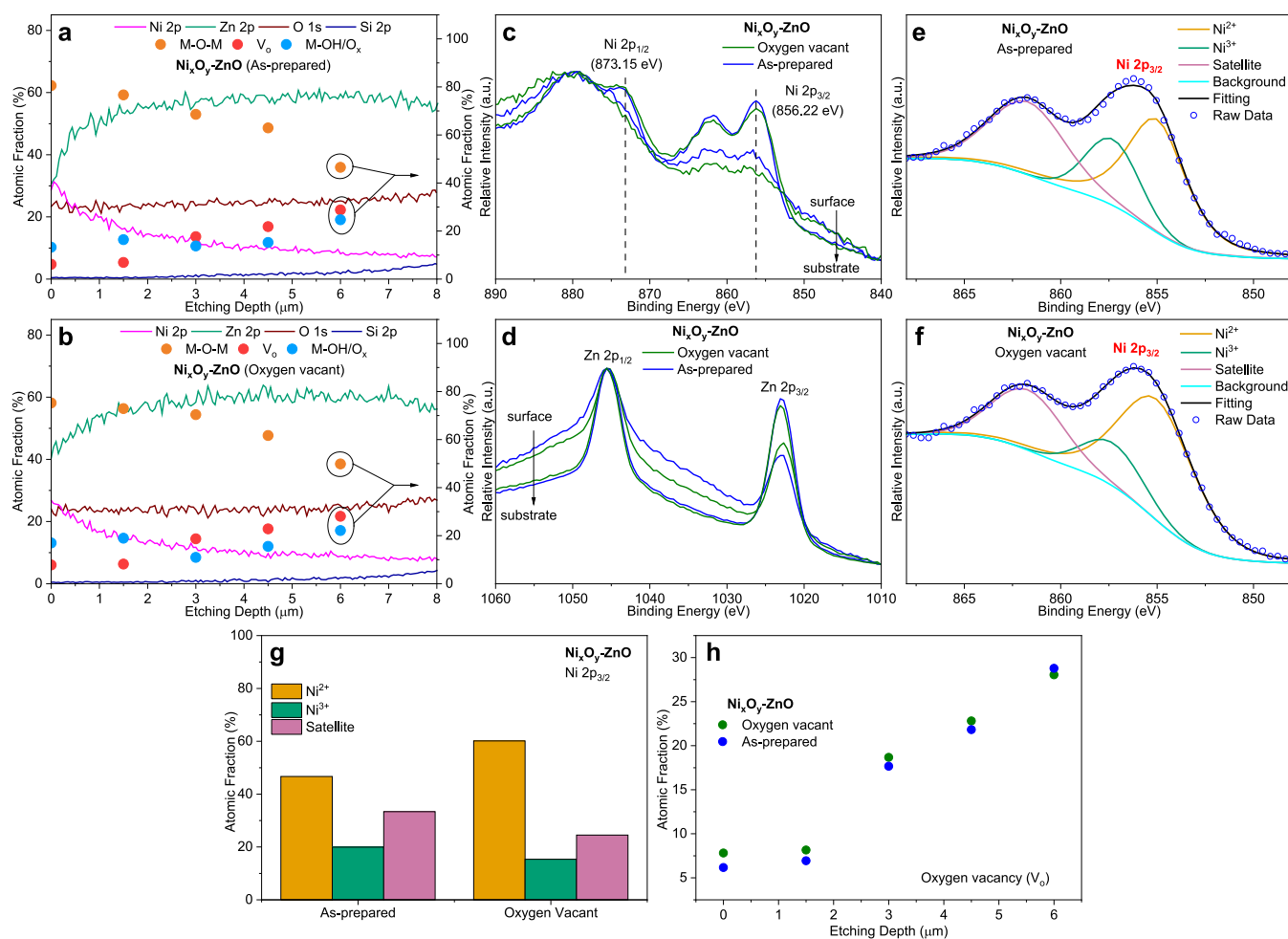


Figure 4. XPS depth profiles as a function of the ion etching depths of (a) as-prepared and (b) oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks showing the fractions of oxygen (brown line), Zn (green line), Ni (pink line), Si (blue line), M-O-M (orange circle), V_o (red circle), and M-OH/ O_x (blue circle) content. High-resolution XPS spectra of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks at various depths showing the (c) Ni 2p and (d) Zn 2p. Deconvoluted Ni 2p spectra of (e) as-prepared and (f) oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks. Atomic fractions of (g) Ni^{2+} and Ni^{3+} and (h) oxygen vacancies in as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks.

and Zn, while no Ni peaks were observed for the pure ZnO layers (Figure 3b). This confirms the successful deposition of Ni_xO_y on the ultraporous ZnO nanoparticle networks. UV-vis analysis was conducted to examine the optical properties of the films after Ni deposition and DUV photoactivation. Figure 3c displays the optical absorbance of the pure ZnO and of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks. The absorbance of all films decreases at a wavelength of ca. 375 nm, which corresponds to the direct band gap of ZnO (3.35 eV). An optical band gap (E_g) of 3.28 eV was estimated by the Tauc plot for all of the nanoparticle networks (Figure S4c). The optical analysis of the heterostructures shows that there was no change in the optical properties of the films after Ni deposition or after the generation of oxygen vacancies *via* DUV photoactivation.

X-ray photoelectron spectroscopy (XPS) was used to analyze the changes in the composition and state of the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks after DUV. Depth profiling was conducted to examine the distribution of oxygen vacancies and the valence state of metal ions throughout the cross-section of the nanoscale heterojunction networks. XPS scans were recorded after layer-by-layer Ar ion

etching. Only Zn, Ni, O, and Si (from the substrate) were detected (see Experimental Section). The absence of other peaks reveals a high purity of the films made by flame-aerosol deposition. All binding energies were corrected by using the adsorbed carbon dioxide peak as a reference in the carbon 1s (C 1s) spectra.⁴⁸ The XPS survey spectra of $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks before and after the DUV photoactivation are shown in Figure S5a. Figure 3d–f present the high-resolution O 1s spectra of the top surface of both the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks. The fraction of oxygen defects increased by 10.5% after DUV photoactivation, and the fraction of loosely bound surface oxygen groups decreased by 26% (Figure 3e, f). A summary of the XPS characteristics in the O 1s of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks is given in Table S1.

The results of the XPS depth profiling are shown in Figure 4a and b before and after DUV photoactivation, respectively. The complete depth profiles of the elements detected in Figure S5b–c show a decrease in metal peaks and an increase in the Si (substrate) peak with increasing etching depth and thus proximity to the Si

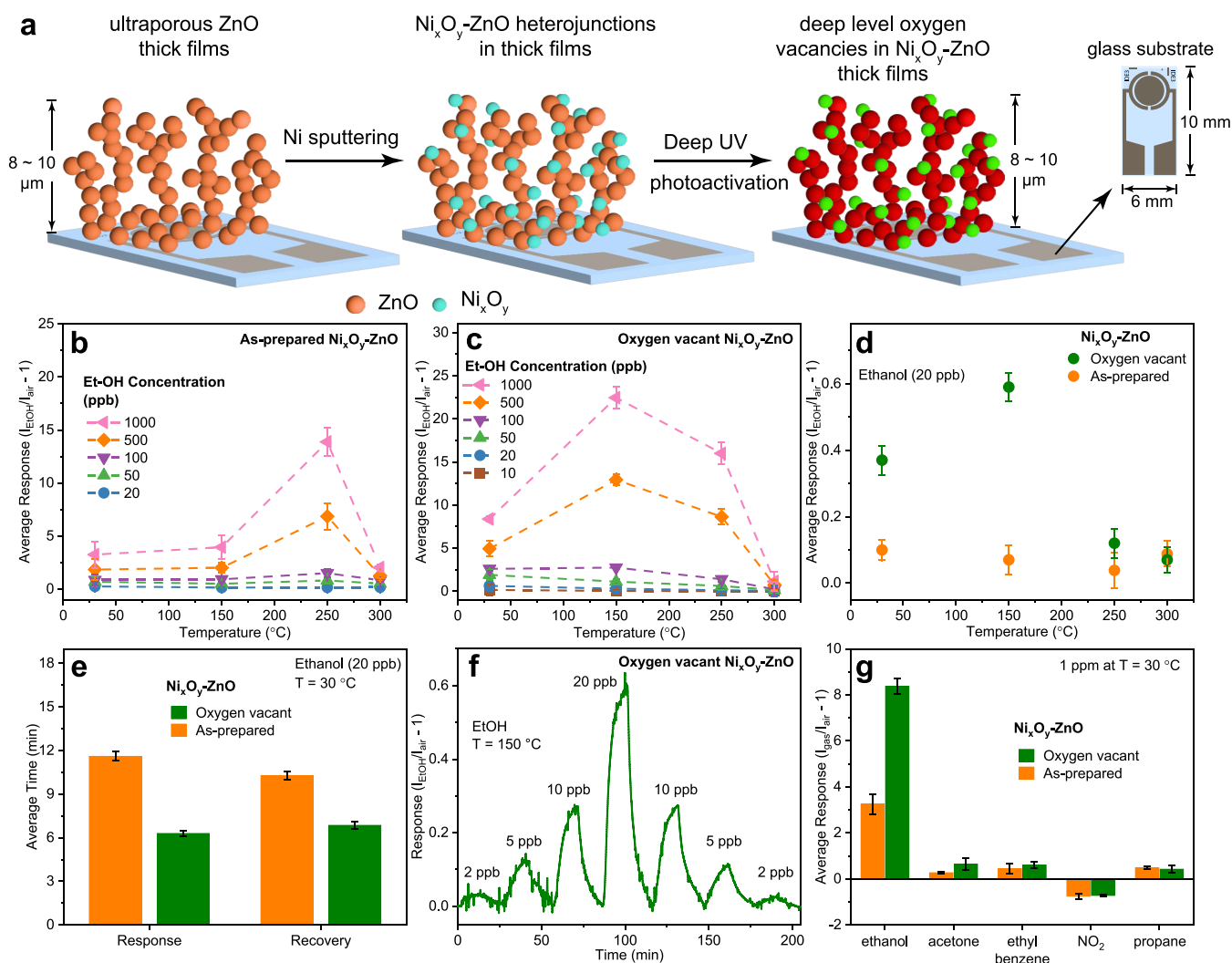


Figure 5. (a) Schematic of the photochemoresistive sensor fabrication. Average ethanol sensor response as a function of concentration and temperature of (b) as-prepared and (c) oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$. (d) The average sensor response to 20 ppb ethanol as a function of temperature. (e) Average response and recovery times of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ to 20 ppb ethanol at room temperature. (f) Representative sensor response dynamic of the oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks to low concentration (2–20 ppb) of ethanol at the optimal temperature of 150 °C. (g) Selectivity of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks to 1 ppm of propane, NO_2 , ethylbenzene, acetone, and ethanol at room temperature.

1s signal intensity is due to the contribution of oxygen from the native oxide of the SiO_2 substrate. The O 1s spectra were deconvoluted into three peaks by a Gaussian–Lorentzian function to determine the fraction of oxygen vacancies. These peaks represent lattice oxygen groups (530.25 ± 0.2 eV, M–O–M), oxygen defects (531.05 ± 0.2 eV, V_O) in ZnO and Ni_xO_y , and loosely bound oxygen groups (532.19 ± 0.2 eV, M–OH/ O_x) such as hydroxyl groups or oxygen molecules adsorbed on the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks (Figure S6).^{3,21,49} Table S2 shows the calculated values of each oxygen species with and without DUV photoactivation at various depths of the nanoscale heterojunction networks. High-resolution spectra of O 1s, Zn 2p, and Ni 2p at all etching levels are provided in Figure S7.

Despite the significant thickness of the ZnO nanoparticle networks and the line-of-sight nature of the sputtering process, the presence of Ni was observed throughout the whole cross-section of the nanoscale heterojunction networks. This is attributed to the highly porous fractal structure of the aerosol self-assembled ZnO nanoparticle networks that allowed Ni to

penetrate deeply into the nanoparticle networks, creating localized *p-n* nanoscale heterojunction networks of $\text{Ni}_x\text{O}_y\text{-ZnO}$ (Figure 4c). The composition of the deposited Ni was also uniform throughout the cross-section, with the Ni 2p spectra showing a consistent split into $2p_{1/2}$ and $2p_{3/2}$ both near the top surface of the network and close to the substrate (Figure 4c).

Figure 4d displays high-resolution Zn 2p spectra of the as-prepared and oxygen vacant nanoscale heterojunction networks. Both spectra show two peaks in Zn 2p spectra due to spin–orbital coupling, which are $2p_{1/2}$ and $2p_{3/2}$, respectively. There was no peak shift in the oxygen vacant nanoscale heterojunction networks, obtained by DUV exposure. This excludes the substitutional replacement of Ni in the Zn sites, which would impact the bonding energy of the local Zn atoms. This was also confirmed by XRD (Figure 3a–c) indicating that the sputtered Ni atoms did not dope in the ZnO crystal lattice but formed segregated domains. The relative increase in the intensity of Zn 2p with increasing etching depth is attributed to the lesser amount of Ni penetrating close to the substrate. The

Ni $2p_{3/2}$ spectra of $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks were deconvoluted into three peaks by using a Gaussian–Lorentzian function. These peaks confirmed the presence of two valence states of Ni: Ni^{2+} around 854.5 eV and Ni^{3+} around 857.8 eV. The deconvoluted Ni $2p_{3/2}$ spectra of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks are shown in Figure 4e, f. The results show that NiO is the dominant nickel oxide phase in the as-prepared nanoparticle networks, while a minor amount of Ni_2O_3 was also present. DUV photoactivation further decreased the fraction of Ni^{3+} (Figure 4g), in line with the induction of oxygen vacancies in the nanoscale heterojunction networks. The fraction of oxygen defects at various depths of the heterostructures increased as a result of UV exposure, as seen in Figure 4h.

To investigate the impact of the oxygen vacancy on the photocatalytic activity and electrical properties of the nanoscale heterojunction networks, chemical sensors were fabricated on commercial Pt interdigitated electrode substrates (Experimental Section) by the same flame-aerosol self-assembly (Figure 5a), with the same ZnO thickness and Ni-loading as shown in the above characterizations. The thickness of the ZnO layer was selected in line with previously best-performing chemical sensors and UV-photodetectors.^{10,11,50} Ethanol was used as a test analyte to interrogate the photochemoresistive properties of the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks before and after the induction of oxygen vacancies. The sensor response was measured as a change in current at a constant voltage of 1 V. All the sensors were constantly exposed to synthetic air and solar illumination (AM1.5 solar simulator, 67 mW cm^{-2}) while cycling the exposure to the analyte gas on/off during 15 min cycles. The sensing experiments were performed three times at temperatures between 30 and 300 °C with ethanol concentrations ranging from 10 to 1000 ppb. The average response was reported with the standard error as an error bar.

The representative dynamic response of the as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks to ethanol concentrations between 10 ppb and 1 ppm are shown in Figure S8a, b. The sensor response increased with an increasing ethanol concentration. Notably, the sensors were able to detect very low concentrations of ethanol at low temperatures (30–150 °C). However, the as-prepared $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks could not detect 10 ppb of ethanol at any temperature, while the oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks achieved the detection of 10 ppb of ethanol already at room temperature (Figure S8a, b). Furthermore, Figure 5b, c shows that the induction of oxygen vacancies decreased the optimal sensor operating temperature from 250 to 150 °C and significantly increased the sensor response. At room temperature, the sensor response to 1 ppm of ethanol increases by 88% and by 140% at 150 °C.

The response to the trace concentration (20 ppb) of ethanol also increases by 83% at room temperature (Figure 5d). Furthermore, the response and recovery times to 20 ppb ethanol at room temperature decrease by 59% and 40% after DUV photoactivation (Figure 5e). Notably, the oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks show a better response at lower temperatures as compared to the as-prepared $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks. However, at high temperatures, the response decreases. This is attributed to the oxygen vacancies becoming more mobile and the nanoscale

heterojunctions becoming more catalytically active, causing the combustion of ethanol molecules before they reach the most sensitive bottom part of the nanoscale heterojunction networks near the interdigitated electrodes as suggested for other highly catalytic active chemoresistive films.⁵¹ To assess the long-term stability of oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks, the sensors were tested after one month of initial experiments at 150 °C. A decrease of 159% was observed for 1 ppm of ethanol at 150 °C (Figure S9a). Furthermore, the ethanol testing was carried out with oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ sensors in realistic conditions for the possible application as breath analyzers in relative humidity of 30 and 50% at room temperature under dark conditions (Figure S9b). The sensors did not respond to any concentration of ethanol gas. This was attributed to the lower number of reaction sites available to the ethanol in the presence of an interfering gas such as H_2O molecules, a low operating temperature of the sensors as well as the absence of solar light.

Overall, the above characterization shows that the generation of oxygen vacancies significantly improves the sensitivity, lowers the detection limit, and improves the response and recovery times of the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks enabling their operation at lower temperatures. Notably, our nanoscale heterojunction networks achieve the highest responsivity reported among recently reported sensors based on nano heterostructures with oxygen defects (Table 1). It should also be noted that, here, the oxygen vacancies in the $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks were effectively generated using our cost-effective low-power (30 mW cm^{-2}) DUV photoactivation process at 200 °C, compared to other expensive, multistep, and high-temperature methods which may damage and/or compromise the stability of the fragile nanoscale heterostructure networks.

$\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks with oxygen vacancies showed a very low limit of detection of 2 ppb of ethanol at 150 °C. Figure 5f shows the dynamic sensor response to extremely low concentrations of ethanol (2–20 ppb) at this optimal temperature (150 °C). Furthermore, poor selectivity at low temperatures is a drawback of metal oxide sensors. Here, the oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks displayed excellent selectivity for ethanol close to room temperature (30 °C) against other volatile organic compounds (VOCs) (Figure 5g). Notably, the ethanol selectivity increased 30 times after DUV photoactivation. The latter led to no significant improvement in the response to propane and NO_2 gases, while an increase in the response toward acetone and ethylbenzene was observed at room temperature. This difference in the response was tentatively attributed to the differences in the polarity and size of these VOC molecules. Generally, ethanol molecules show higher binding affinities as the bipolar hydroxyl group (OH) interacts with the surface *via* two different charge transfer mechanisms, unlike monopolar molecules like acetone. This has been reported recently for SnO_2 at room temperature and confirmed both experimentally and theoretically.⁵² Notably, both as-prepared and oxygen vacant $\text{Ni}_x\text{O}_y\text{-ZnO}$ nanoscale heterojunction networks also showed an inverse response to NO_2 gas as expected for oxidizing analyte and *p*-type semiconductors (Figure 5g).

To confirm this mechanism and gain further insights into the role of the oxygen vacancy in enhancing the response and selectivity to ethanol, DFT calculations were carried out to study the effects of surface adsorption and decomposition of

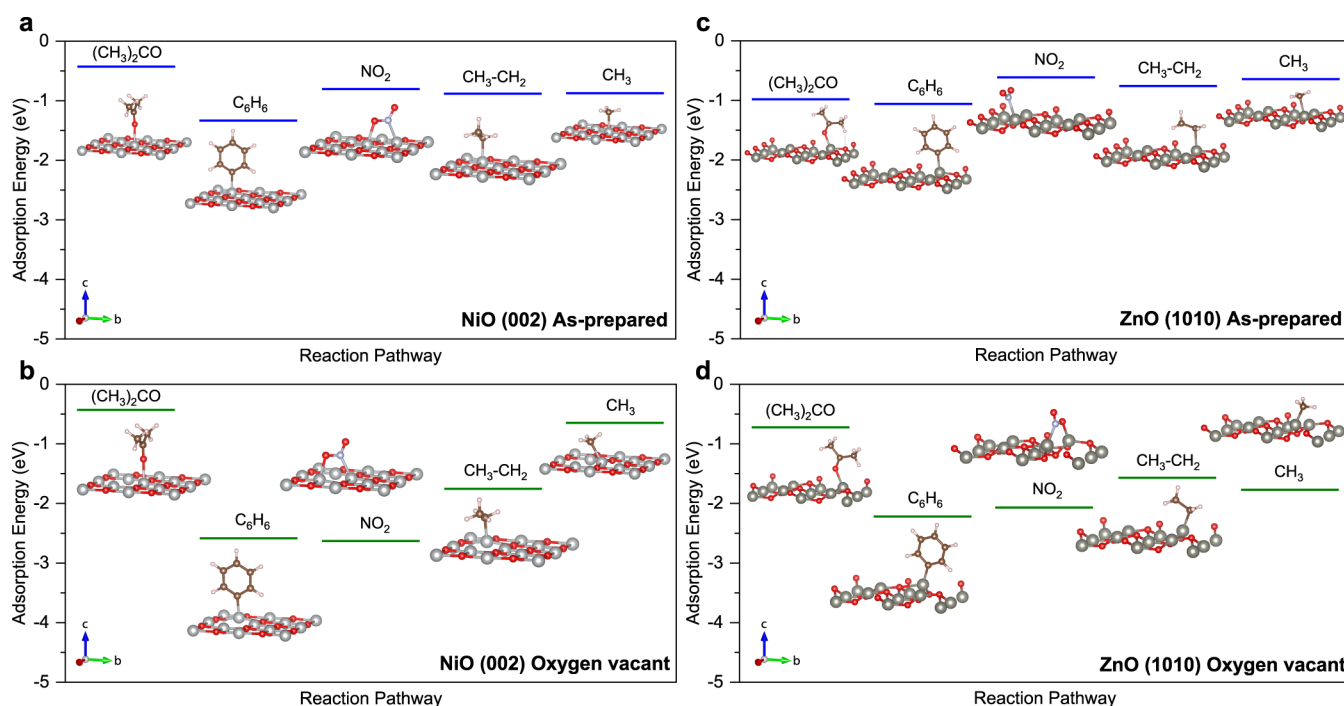


Figure 6. Adsorption energies of various VOCs with corresponding geometric configurations calculated by theoretical analysis using DFT of NiO (a) as-prepared and (b) oxygen vacant and of ZnO (c) as-prepared and (d) oxygen vacant.

VOCs on the *p*-type NiO and *n*-type ZnO surfaces with and without DUV photoactivation. The preferential planes of (002) and (1010) were selected for NiO and ZnO nanoparticle networks with cubic and hexagonal crystal structures, respectively. Geometrical optimizations were achieved by relaxing all ionic positions and supercell vectors until the Hellman-Feynmann forces were less than 0.02 eV^{-1} . The dispersion correction was also considered in this study by using the DFT-D3 method.⁵³ The slab method was used to model NiO and ZnO surfaces with surface-bulk-like interactions by fixing the coordinates of the atoms in three bottom layers and relaxing the coordinates of the atoms in the three uppermost layers. The results show that all VOC molecules dissociate at the surface of the nanoparticles and that the corresponding adsorption energies increase in the presence of oxygen vacancies (Figure 6a–d).

The dissociation of VOC molecules can occur through diverse pathways that involve intricate reaction kinetics. In particular, ethylbenzene (C_8H_{10}) molecules can undergo direct reactions with the sensing layer to produce a response signal, or they can decompose into C_6H_6 , C_3H_8 , or CH_3 . The adsorption energy of NO_2 molecules on *p*-type NiO increased from 0.8 to 2.6 eV, while that on *n*-ZnO rose from 0.6 to 2.1 eV after the introduction of oxygen vacancies. This resulted in a cumulative increase of 108% in the adsorption energy of NO_2 molecules in Ni_xO_y -ZnO nanoscale heterojunction networks due to oxygen vacancies. Notably, the rise in the adsorption energies of NO_2 molecules was more significant in *n*-type ZnO compared to *p*-type NiO. However, the increase was highest in the case of ethanol molecules. Furthermore, this increase in ethanol adsorption energy was higher ($\sim 98\%$) in the *p*-type NiO as compared to *n*-type ZnO ($\sim 59\%$). This indicates that the formation of *p*-*n* Ni_xO_y -ZnO nanoscale heterojunction networks and the presence of oxygen vacancies in both metal oxides can be used to enhance the selectivity of metal oxide gas sensors by surface energy engineering. The adsorption energies

for different VOCs on both as-prepared and oxygen vacant Ni_xO_y -ZnO nanoscale heterojunction networks are summarized in Table 2.

Table 2. Summary of Adsorption Energies of Various VOCs on As-Prepared and Oxygen Vacant Ni_xO_y -ZnO Nanoscale Heterojunction Networks

VOC Species	As-prepared NiO(002)	Oxygen vacant NiO(002)	As-prepared ZnO(1010)	Oxygen vacant ZnO(1010)
$(\text{CH}_3)_2\text{CO}$ (acetone)	−0.43	−0.42	−0.99	−0.74
$\text{C}_2\text{H}_5\text{O}$ (from $\text{C}_2\text{H}_5\text{OH}$ -ethanol)	−1.27	−3.69	−1.51	−2.78
Benzene (from C_8H_{10} -ethylbenzene)	−1.35	−2.59	−1.07	−2.23
NO_2 (Nitrogen dioxide)	−0.82	−2.64	−0.60	−2.07
$\text{CH}_3\text{-CH}_2$ (from C_3H_8 -propane or ethylbenzene)	−0.91	−1.75	−0.75	−1.57
CH_3 (from propane, acetone, ethylbenzene, ethanol)	−0.89	−2.00	−0.65	−1.78

Given the above experimental characterization and modeling results, here, we attribute the strong increase in the photoconductive sensor performance to a higher photocatalytic activity of the surface due to the formation of oxygen vacancies in both Ni_xO_y and ZnO and the localized *p*-*n* nanoscale heterojunctions between Ni_xO_y and ZnO that facilitate the photoexcited charge separation making them available for the interaction with the analyte molecules. Furthermore, the ultraporous morphology of the electron-depleted ($d_p < \text{Debye length}$) Ni_xO_y -ZnO nanoscale heterojunction networks facilitate the diffusion of the analyte gas and products in and out of the sensing element resulting in

an increased active area. At ambient conditions, oxygen adsorbs on the surfaces of Ni_xO_y and ZnO and traps electrons, forming O_2^- , and increasing the depletion region^{19,24} (Figure S10 and 1a). The mismatch between the bandgap and Fermi level causes electrons to flow from Ni_xO_y to ZnO and holes flow in the opposite direction at Ni_xO_y -ZnO nanoscale heterojunctions, effectively separating electrons, and holes (Figure 1b). The photocharge separation at the Ni_xO_y and ZnO also allows for charges to be available on the ZnO and Ni_xO_y surface leading to, respectively, more and less oxygen chemisorption, increasing the number of reaction sites available for target gas molecule, such as ethanol on the surface of the nanoscale heterostructures.

CONCLUSION

In summary, we have reported a flexible approach for the engineering of the surface activity of high-performance three-dimensional nanoscale heterojunction networks, demonstrating vastly enhanced photoactivity and sensing performance. Ultraporous nanoparticle networks of Ni_xO_y -ZnO nanoscale heterojunctions were self-assembled by flame-aerosol deposition and sputtering leading to highly uniform, electron-depleted porous layers. Oxygen vacancies were induced throughout the whole Ni_xO_y -ZnO nanoscale heterojunction networks in a controllable manner by a moderate temperature deep UV photoactivation. XPS depth profiling showed the presence of localized p - n nanoscale heterojunctions and oxygen vacancies throughout the Ni_xO_y -ZnO whole 8–10 μm thick networks. Upon induction of oxygen vacancies, the sensing performance of the nanoscale heterojunction networks increased substantially. The sensitivity to ethanol, as a representative analyte, was enhanced by 88% with a detection limit of 2 ppb at 150 °C and a 30-fold increase in selectivity against a panel of other VOCs at room temperature. Theoretical study of the as-prepared and oxygen vacant nanoscale heterojunction surface by DFT revealed that the adsorption energy of ethanol increased by 80% in the oxygen vacant Ni_xO_y -ZnO nanoscale heterojunction networks, explaining the strong increase in photochemoresistive response of the DUV photoactivated devices. Overall, the high specific surface area, efficient separation of electron–hole pairs at the nanoheterojunction interfaces, deep penetration of target gas molecules into the porous networks, and oxygen-vacant surface all contributed to the outstanding performance of this nanoscale architecture, achieving record-high sensitivity and selectivity at room and moderated temperatures. We believe that these findings provide a facile approach for the engineering of highly performing photocatalytic materials with applications including miniaturized sensing, photocatalysis, and a variety of optoelectronic applications.

EXPERIMENTAL SECTION

Materials. All materials were used as received without any further purification. Zinc naphthenate (10% Zn, Merck) was used as a precursor and dissolved in xylene (Merck) to generate a solution with a total Zn metal atom concentration of 0.3 M. The solution was stirred for 30 min at room temperature to ensure complete dispersion.

Fabrication of Ni_xO_y -ZnO. Flame spray pyrolysis (FSP) was used to synthesize and for direct deposition of ZnO nanoparticles on Pt-interdigitated electrodes.^{44,45} The Zn precursor solution was delivered via an automatic syringe pump (New Era syringe pump; Model: NE-300) at a flow rate of 5 mL min^{-1} and dispersed into a fine spray with 7 L min^{-1} oxygen at a constant pressure drop of 4.5 bar. The spray was ignited by supporting premixed methane/oxygen flames ($\text{CH}_2 =$

1.8 L min^{-1} , and $\text{O}_2 = 2$ L min^{-1}). Glass substrates containing 180 pairs of Pt interdigitated electrodes of 5 μm width and spacing (Micrux, Spain) were placed on a water-cooled substrate holder at a height above burner (HAB) of 20 cm. All substrates were cleaned three times with ethanol before deposition. After FSP deposition, the substrates were heat-treated at 300 °C for 12 h in a furnace (Brother High-Temperature Furnace XD-1.2KN) to stabilize the nanoparticle size and avoid resistive sintering during the gas sensing measurements. Ni was deposited by an AJA ATC 2400 magnetron sputtering system in an ambient Ar/ O_2 chamber (18/2 sccm). The chamber was evacuated to 9×10^{-7} Torr before deposition and sputtering deposition was performed at a gas pressure of 4×10^{-3} Torr for 45 s at room temperature with a power of 200 W. The equivalent of a 6 nm Ni layer bulk thickness was sputtered onto these ultraporous ZnO nanoparticle networks and thereafter sintered at 300 °C for 12 h to form localized Ni_xO_y -ZnO p - n nanoscale heterojunctions and improve interparticle necking and connectivity.

For the DUV photoactivated modulation of FSP-synthesized ZnO films, a customized temperature-controllable system was adopted for precise temperature control, and a homemade DUV irradiation system was designed by using an excimer lamp (UV emission wavelength of 172 nm, 20×12 cm², Wonik Q & C) with the constant intensity of UV ~ 30 mW cm^{-2} . Dry N_2 gas was supplied to the customized chamber to prevent the formation of ozone (O_3) and the films and sensor devices were identically heated at 200 °C for 2 h under UV.

Material Characterization. Topographical and optical characterization techniques were used to analyze the morphology of the samples. Analytical scanning electron microscopy, (Zeiss Ultra plus field emission scanning electron microscope (FESEM)) was used for morphological analysis at 3 kV. The results were analyzed using ImageJ software and the FracLac plugin to calculate the fractal dimensions and lacunarity.^{45,54} This algorithm covers the SEM image of the analyzed object with a grid of decreasing dimensionality and counts the number of boxes required to fully capture the solid image. In this way, the structure of the pattern is captured with a high accuracy. Information about the morphology, particle size, and lattice plane spacing was acquired by using a high-resolution transmission electron microscope (JEOL 2100F HRTEM) operated at 200 kV. A small amount of powder sample was added to ethanol and sonicated for 20 min to achieve a stable suspension, followed by casting a drop onto Lacey carbon 200 mesh nickel grids. The particle size distribution was evaluated by manually measuring the size of 625 nanoparticles using ImageJ. The simulation of the polycrystalline TEM diffraction patterns was performed using the JEMS⁵⁵ code as implemented on version 4.6131U2018. A Bruker D2 Phaser diffractometer was utilized to perform X-ray diffraction (XRD). Each sample was scanned using a Cu $K\alpha$ (1.54) radiation source with an operating voltage of 30 kV and a current of 10 mA. The XRD patterns were recorded with a scan rate of 0.75° min^{-1} in the 2θ range of 10–80° at an increment of 0.02°. The modified Scherrer equation⁴⁷ was applied to accurately calculate the crystalline domain size. XRD patterns were analyzed by commercially available software HighScore Plus using the Inorganic Crystal Structure Database (ICSD). The modified Scherrer equation is used to accurately estimate the size of nanoparticles from XRD peaks⁴⁷ and is given by

$$\ln \beta = \ln \frac{1}{\cos \theta} + \ln \frac{K\lambda}{D} \quad (1)$$

where β is full width at half-maximum, in radians, located at any 2θ position in the pattern, λ (nm) is the radiation of wavelength, K is the Scherrer constant or shape factor, and D is the particle size in nm.

XRD data were recorded at room temperature on a PANalytical Empyrean Series 1 diffractometer fitted with a 1D PIXcel detector. Standard Bragg–Brentano data collections utilized fixed diversion slits in the incident beam and a Ni filter to remove Cu $K\beta$ radiation to achieve Cu $K\alpha_1$ and Cu $K\alpha_2$ radiation source. Grazing incident experiments (GI-XRD) utilized a hybrid monochromator to achieve a parallel Cu $K\alpha$ radiation source; the incident angle is reported with the data.

UV–Vis transmittance and absorption spectra were measured using a PerkinElmer (Lambda 1050 UV/vis/NIR) spectrophotometer. The characterization for both atomic orbital and chemical composition was analyzed by X-ray photoelectron spectroscopy (XPS, K-APLPHA+, Thermo Fisher Scientific) using a monochromated Al–K α source. Survey and high-resolution spectra were obtained with pass energies of 100 and 20 eV. Further depth profiling was carried out through a repetitive process of Ar etching and an immediate XPS scan. The ultraporous ZnO films were deposited on SiO₂/Si substrates for SEM, XRD, and XPS analysis while on quartz substrates for UV–vis analysis.

The porosity, ϵ , was estimated from gravimetric and cross-sectional SEM analysis as previously suggested⁵⁶

$$\epsilon = 1 - \frac{T_{\text{bulk}}}{T_{\text{SEM}}} \quad (2)$$

where T_{SEM} is the thickness of the ZnO fractals measured from cross-sectional SEM, while T_{bulk} is the equivalent bulk thickness given by

$$T_{\text{bulk}} = 1 - \frac{V_{\text{ZnO}}}{A} = \frac{m}{\rho_{\text{ZnO}}A} \quad (3)$$

here, m is the mass of deposited material onto the substrate of the area A and ρ is the average density of the fractal media.

Sensing Measurements. The sensing measurements were performed as follows: VOCs (ethanol, 9.91 ppm in N₂ (BOC Ltd.); acetone, 10.1 ppm in N₂ (BOC Ltd.); ethylbenzene, 10 ppm in N₂ (BOC Ltd.); and propane, 10 ppm in N₂ (BOC Ltd.)) were controlled by mass flow controllers (Bronkhorst) and further diluted with simulated air (0.1 L min⁻¹ O₂ + 0.4 L min⁻¹ N₂ (BOC Ltd.)) to reach the desired concentration (from 2 ppb to 1 ppm), and the total gas flow rate was still kept at 0.5 L min⁻¹. The sensors were heated for 2 h at 300 °C in the continuous flow of synthetic air to desorb any atmospheric contamination before the start of each sensing experiment. The temperature of a stage in the sensing chamber (Linkam) was controlled by a temperature controller and the sample was illuminated through a quartz window by a solar simulator (NewSpec, LCS-100) with an AM1.5 filter glass. The dynamic response of the gas sensor was recorded by an electrochemical workstation (CHI 660E, USA) under a constant applied voltage of 1 V. Standard errors were calculated in the sensing measurements by performing sensing experiments at least 3 times with the same experimental conditions and the average responses are presented as the sensor response. The sensing response of the sensors was calculated as

$$\text{Sensor Response} = \left(\frac{I_{\text{gas}}}{I_{\text{air}}} \right) - 1 \quad (4)$$

where I_{gas} is the current in the presence of the target gas and I_{air} is the current in the air without the target gas. The target gas flowed into the sensing chamber for 15 min and was then switched off for 15 min. The response and recovery times were calculated using a threshold of 90% change in the response in the presence and absence of the target gas, respectively.

DFT Calculations. Spin-polarized DFT calculations were implemented using the Vienna ab initio simulation package (VASP)⁵⁷ with the core and valence electronic interactions being modeled using the projector augmented wave (PAW) method.⁵⁸ The Perdew–Burke–Ernzerhof (PBE) exchange–correlation function was employed.⁵⁹ The wave function was expanded with a kinetic energy cutoff of 500 eV and Gamma k-points were used. The Dudarev approach was implemented for the running of the DFT+ U calculation, considering the electron self-interaction error in Ni and Zn. Hubbard's parameters of $U = 5.77$ for Ni in NiO,⁶⁰ $U = 10.0$, and $U = 7.0$ for Zn and O in ZnO were used.^{61,62} Geometrical optimizations were achieved by relaxing all ionic positions and supercell vectors until the Hellman–Feynman forces were less than 0.02 eV⁻¹. The dispersion correction was also considered in this study by using the DFT-D3 method.⁵³ The slab method was used to model NiO and ZnO surfaces with surface–bulk-like interactions, by fixing

the coordinates of the atoms in three bottom layers and relaxing the coordinates of the atoms in the three uppermost layers. The Gamma k-points of $3 \times 3 \times 1$ and $5 \times 3 \times 1$ were used for sampling the Brillouin zone for (2×2) NiO (002) and (2×2) ZnO (10 $\bar{1}0$) surfaces. A vacuum region of 20 Å was introduced in the direction of the z -axis to avoid interactions between the periodic images. The adsorption energies (E_{ads}) are computed using the following equation:

$$E_{\text{ads}} = E_{(\text{slab}+\text{adsorbates})} - E_{\text{slab}} - E_{\text{adsorbates}} \quad (5)$$

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsami.3c12668>.

Flame spray schematic, proposed reaction pathways and surface decomposition of ethanol on Ni_xO_y-ZnO nanoscale heterojunction networks, top-view and side-view SEM images, fractal analysis, XRD of pure NiO, NiO-ZnO, Tauc plots of pure ZnO, as-prepared and oxygen vacant NiO-ZnO, XPS survey and complete XPS depth profiles of Ni_xO_y-ZnO p - n nanoscale heterojunction networks; deconvoluted O 1s at various etching depths of as-prepared and oxygen vacant Ni_xO_y-ZnO; complete high resolution XPS spectra of O 1s, Zn 2p, and Ni 2p of as-prepared and oxygen vacant Ni_xO_y-ZnO p - n nanoscale heterojunction networks, electronic band structure of Ni₂O₃, gas sensing mechanism for Ni_xO_y-ZnO p - n nanoscale heterojunctions, stability and humidity tests, XPS characteristic summary of O 1s at the top surface of as-prepared and oxygen vacant Ni_xO_y-ZnO p - n nanoscale heterojunctions, Fractions of oxygen species in O 1s of as-prepared and oxygen vacant Ni_xO_y-ZnO fractals at various etching depths (PDF)

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Notes

The authors declare no competing financial interest.

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